

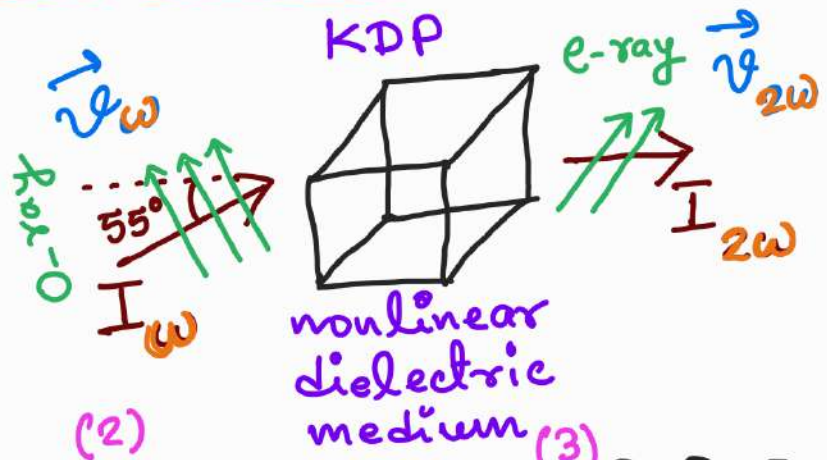
# NON LINEAR OPTICS

polarization ↑

free space permittivity ↑

$$P = \epsilon_0 \chi E$$

↓ electric field  
 ↓ electric polarizability



$$P_i = \epsilon_0 \left[ \chi_{ij}^{(1)} E_j + \chi_{ijk}^{(2)} E_j E_k + \chi_{ijkl}^{(3)} E_j E_k E_l + \dots \right]$$

(1) 9 component ≡ 6 independent component  
 (2) 27 component ≡ 18 independent component  
 (3) 81 component ≡ 54 independent component

# isotropic sample:  $P = \epsilon_0 (\chi^{(1)} E + \chi^{(2)} E^2 + \chi^{(3)} E^3 + \dots)$

$$E = E_0 \sin \omega t \quad (\text{say}) \Rightarrow \epsilon_0 \chi^{(1)} E_0 \sin \omega t + \epsilon_0 \chi^{(2)} E_0^2 \sin^2 \omega t + \epsilon_0 \chi^{(3)} E_0^3 \sin^3 \omega t + \dots$$

$$= \epsilon_0 \chi^{(1)} E_0 \sin \omega t + \frac{\epsilon_0 \chi^{(2)}}{2} E_0^2 (1 - \cos 2\omega t) + \frac{\epsilon_0 \chi^{(3)}}{4} E_0^3 (3 \sin \omega t - \sin 3\omega t)$$

DC polarization (optical rectification)

frequency doubling / 2nd harmonic generation (SHG)

## NO SHG

- ① Has center of symmetry / inversion center (calcite)
- ② Not piezoelectric. quartz, KDP, ADP ⇒ Piezoelectric

$$I_{2\omega} \propto \sin^2 \left[ \frac{2\pi (n_\omega - n_{2\omega}) l}{\lambda_0} \right] \Rightarrow l_c = \frac{\lambda_0}{4(n_\omega - n_{2\omega})}$$

maximum  $I_{2\omega}$

2w

$$(n_w - n_{2w})^2$$

coherence length

Index Matching:  $n_w = n_{2w}$

Frequency mixing:  $E = E_{01} \sin \omega_1 t + E_{02} \sin \omega_2 t$

$$\Rightarrow P = \chi^{(1)} \text{ term} + \epsilon_0 \chi^{(2)} (E_{01}^2 \sin^2 \omega_1 t + E_{02}^2 \sin^2 \omega_2 t + 2E_{01} E_{02} \sin \omega_1 t \sin \omega_2 t) + \dots$$

$$= \chi^{(1)} \text{ term} + \frac{\epsilon_0 \chi^{(2)}}{2} E_{01}^2 (1 - \cos 2\omega_1 t) +$$

difference photon

$$\frac{\epsilon_0 \chi^{(2)}}{2} E_{02}^2 (1 - \cos 2\omega_2 t) + \epsilon_0 \chi^{(2)} [\cos(\omega_1 - \omega_2)t - \cos(\omega_1 + \omega_2)t]$$

sum photon

Maxwell's eq<sup>n</sup> for linear, homogeneous, isotropic medium:

- ①  $\nabla \cdot \bar{D} = \rho$   $\xleftrightarrow{\text{Gauss's law}}$   $\oint_S \bar{D} \cdot d\bar{S} = \iiint_V \rho dV$
- ②  $\nabla \cdot \bar{B} = 0$   $\xleftrightarrow{\quad}$   $\oint_S \bar{B} \cdot d\bar{S} = 0$
- ③  $\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t}$   $\xleftrightarrow{\text{Faraday's law / Lenz's law}}$   $\oint_C \bar{E} \cdot d\bar{l} = -\iint_S \frac{\partial \bar{B}}{\partial t} \cdot d\bar{S}$
- ④  $\nabla \times \bar{H} = \bar{J}_c + \frac{\partial \bar{D}}{\partial t}$   $\xleftrightarrow{\text{Ampere's law}}$   $\oint_C \bar{H} \cdot d\bar{l} = \iint_S (\underbrace{\bar{J}_c}_{\text{Ampere}} + \underbrace{\frac{\partial \bar{D}}{\partial t}}_{\text{Maxwell}}) \cdot d\bar{S}$

$$\nabla \cdot \bar{J}_c + \frac{\partial \rho}{\partial t} = 0 \rightarrow \text{Continuity}$$

$$\# \nabla \cdot \textcircled{4} \Rightarrow 0 = \nabla \cdot \bar{J}_c + \frac{\partial}{\partial t} \nabla \cdot \bar{D} = -\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial t} \nabla \cdot \bar{D}$$

$\nabla \cdot \bar{D} = \rho - \rho$  where  $\epsilon \neq \epsilon(t)$  But

$C_1 = 0$  otherwise  $\oint_S \bar{D} \cdot d\bar{S} = \iiint_V \rho dV + C_1 V$ , so  $C_1 V$  acts as source which cannot be true.  $\therefore C_1 = 0$

#  $\nabla \cdot \textcircled{3} \Rightarrow 0 = -\frac{\partial}{\partial t} \nabla \cdot \bar{B} \therefore \nabla \cdot \bar{B} = C_2$  where  $C_2 \neq C_2(t)$  but then  $\oint_S \bar{B} \cdot d\bar{S} = C_2 V$  acts as source which cannot be true  $\therefore C_2 = 0$ .

Basically  $\textcircled{1}$  &  $\textcircled{2}$  are subset of  $\textcircled{3}$  &  $\textcircled{4}$ , so 6 scalar equations for 12 unknowns ( $\bar{E}, \bar{B}, \bar{D}, \bar{H}$ ), so 6 more equations (constitutive equations  $\bar{D} = \epsilon \bar{E}, \bar{B} = \mu \bar{H}$ ) are required.

$$\bar{D} = \epsilon \bar{E} = \epsilon_0 \bar{E} + \bar{P} = \epsilon_0 (1 + \chi) \bar{E} = \epsilon_0 \epsilon_r \bar{E}$$

relative  
permeability

$$\# \nabla \times \textcircled{3} \Rightarrow \nabla \times \nabla \times \bar{E} = -\frac{\partial}{\partial t} \nabla \times \bar{B} = -\mu \frac{\partial}{\partial t} \nabla \times \bar{H}$$

$$\therefore \nabla(\nabla \cdot \bar{E}) - \nabla^2 \bar{E} = -\mu \left[ \frac{\partial \bar{J}_c}{\partial t} + \frac{\partial^2 \bar{D}}{\partial t^2} \right] \quad (\because \bar{B} = \mu \bar{H})$$

$$\therefore \nabla^2 \bar{E} = \mu \left[ \frac{\partial \bar{J}_c}{\partial t} + \epsilon \frac{\partial^2 \bar{E}}{\partial t^2} \right] + \nabla \left( \frac{\rho}{\epsilon} \right) \quad (\because \nabla \cdot \bar{E} = \frac{\rho}{\epsilon})$$

$$\therefore \nabla^2 \bar{E} = \mu \frac{\partial \bar{J}_c}{\partial t} + \mu \epsilon \frac{\partial^2 \bar{E}}{\partial t^2} + \frac{1}{\epsilon} \nabla \rho$$

$$\therefore \nabla^2 \bar{E} \approx \mu \epsilon \frac{\partial^2 \bar{E}}{\partial t^2} \quad (\text{In free space } \bar{J}_c = \rho = 0)$$

Similarly  $\nabla^2 \vec{H} = \mu \epsilon \frac{\partial^2 \vec{H}}{\partial t^2} = v^2 \frac{\partial^2 \vec{H}}{\partial t^2}$ ,  $v = \frac{1}{\sqrt{\mu \epsilon}}$

Wave equation solution in Cartesian, cylindrical (Bessel's fn), Spherical polar. coordinate system.

Anisotropic medium:  $D_i = \epsilon_{ij} E_j$ ;  $\epsilon_{ij} = \epsilon_{ji}$

